Sulphur storage in cold molecular clouds: the case of the NH⁺ ⁴ SH[−] salt on interstellar dust grains

J. Vitorino[®][,](http://orcid.org/0000-0001-9676-2605)^{[1](http://orcid.org/0000-0001-6981-0421)★} J.-C. Loison[®],² V. Wakelam®,³ E. Congiu[®]¹ and F. Dulieu®^{1★}

¹*CY Cergy Paris Universite,´ Observatoire de Paris, PSL Research University, CNRS, LERMA, F-95000 Cergy, France* ²*Universite´ de Bordeaux, Institut des Sciences Moleculaires, ´ UMR 5255, F-33400 Talence, France*

³Laboratoire d'Astrophysique de Bordeaux, Université de Bordeaux, CNRS, B18N, allée Geoffroy Saint-Hilaire, F-33615 Pessac, France

Accepted 2024 July 11. Received 2024 July 2; in original form 2024 May 27

ABSTRACT

In comets and in the cold phase of the interstellar medium (ISM), ammonium salts are key molecular species due to their role in the retention of volatile compounds on cold surfaces. In the case of sulphur, the H_2S/OCS ratio observed in protostars could be explained by the presence of ammonium hydrosulphide (NH4SH) salts. However, laboratory data on the properties of NH4SH in ISM cold relevant conditions are rather scarce, as they usually focus on the atmosphere of Jupiter. We propose to consolidate the laboratory data regarding NH4SH on grains, by performing temperature programmed desorption experiments and Fourier transform infrared reflection spectroscopy. The salt was also exposed to H atomsto mimic the ISM conditions. NH4SH wasfound to form *in situ* at 10 K, from a mixture of ammonia (NH₃) and hydrogen sulphide (H₂S). The NH₄⁺ infrared feature (1485 cm⁻¹) is the most prominent one at 80 K. As pure species, H₂S and NH₃ desorb at 76 and 90 K, respectively, whereas they are released into the gas phase at 153 K when adsorbed in the form of salt. The presence of water delays the desorption of the salt until the very end of the water desorption, but does not affect the desorption kinetics. During H-exposure, the salt is dissociated and no new product was detected. As a comparative study, salts have been included in the Nautilus gas–grain model. The results show a good correlation with the observations of IRAS 16293−2422B, as opposed to when NH4SH is not included in the model.

Key words: astrochemistry – molecular processes – methods: laboratory: solid state – ISM: clouds – ISM: molecules.

1 INTRODUCTION

Laboratory data about ammonium hydrosulfide (NH4 ⁺SH−, hereafter NH4SH) characteristics in interstellar medium (ISM) relevant conditions are rather scarce. Indeed, in terrestrial conditions, ammonium hydrosulfide is known to be formed via the mixing of gaseous hydrogen sulfide (H_2S) and ammonia (NH_3) , which are corrosive chemical compounds with high toxicity for humans. Additionally, their use asks for a very precise leak control, especially for H2S, given the intense odour detectable by smell well before the recommended exposure limit concentration.¹ NH₄SH is one of the most interesting salts potentially present in the ISM and in the upper layer of the coma in comets such as 67P/Churyumov– Gerasimenko (Altwegg et al. [2020\)](#page-9-0). The detection of sulphur and ammonia clouds in overlapping layers by the Galileo probe on Jupiter, with a uniform mixture at 10 bars, also indicates that NH_4SH is a major compound of gas-giant planets atmosphere (Niemann et al. [1998;](#page-10-0) Pater et al. [2001;](#page-10-0) Wong et al. [2004;](#page-10-0) Li et al. [2020\)](#page-10-0). This has been supported by theoretical models (Weidenschilling & Lewis [1973;](#page-10-0) Atreya et al. [1999;](#page-9-0) Friedson [2005\)](#page-10-0), as well as spacecraft observations (Roman, Banfield & Gierasch [2013\)](#page-10-0). Consequently, most laboratory

studiesinvestigate the radiolysis and photodissociation of ammonium hydrosulfide. For example, the infrared and radiolytic data produced by Loeffler et al. [\(2015\)](#page-10-0), allow them to predict that NH4SH would be detectable mostly in crystalline form at 200 K in gas-giant clouds, but that ionizing radiations would destroy the crystalline structure, producing an amorphous mixture instead. The first documented laboratory experiments of ultraviolet (UV)-photolysis of mixed H2S, NH₃, and H₂O ices and pure NH₄SH were done by Lebofsky & Fegley [\(1976\)](#page-10-0), in order to compare the spectra obtained with the observations of Jupiter and Saturn satellites, as well as Saturn rings. More recently, combined studies of UV–vis reflectance spectra of photolyzed NH4SH, together with visible light imaging of these mixtures, have proven their usefulness for remote sensing campaigns on Jupiter, showing a colour-dependence on the irradiation dose and temperature (Loeffler et al. [2016;](#page-10-0) Loeffler & Hudson [2018\)](#page-10-0).

Salts are ionic compounds, often formed by proton transfer in an acid–base reaction. The latter can take place at cryogenic temperatures due to a very low-activation energy (Theulé et al. [2013\)](#page-10-0), while the salt created is usually much more refractory than any of the individual species involved in the ionic bond. For instance, NH3 and HCOOH or $CH₃COOH$, when interacting in form of ammonium formate $(NH_4$ ⁺COO⁻) or ammonium acetate $(NH_4$ ⁺CH₃COO⁻), can remain on dust grains far beyond their respective desorption temperature (Bergner et al. [2016;](#page-9-0) Kruczkiewicz et al. [2021\)](#page-10-0). The same goes for reactions such as $NH_3+HCN \rightarrow NH_4$ ⁺CN⁻ (ammonium cyanide, Clutter & Thompson [1969;](#page-9-0) Gerakines, Moore &

© 2024 The Author(s).

Published by Oxford University Press on behalf of Royal Astronomical Society. This is an Open Access article distributed under the terms of the Creative Commons Attribution License [\(https://creativecommons.org/licenses/by/4.0/\)](https://creativecommons.org/licenses/by/4.0/), which permits unrestricted reuse, distribution, and reproduction in any medium, provided the original work is properly cited.

⁻ E-mail: julie.vitorino@cyu.fr (JV); francois.dulieu@cyu.fr (FD)

¹Hydrogen sulfide; SDS No. 295 442 [Print]; Merck Life Science S.A.S.: Saint-Quentin-Fallavier, France, 2024 Jan 3. [https://www.sigmaaldrich.com/](https://www.sigmaaldrich.com/FR/fr/product/aldrich/295442) FR/fr/product/aldrich/295442 (accessed 2024 Mar 4).

Hudson [2004;](#page-10-0) Noble et al. [2013;](#page-10-0) Gerakines, Yarnall & Hudson [2024\)](#page-10-0), or NH_3 + HNCO $\rightarrow NH_4$ ⁺OCN⁻ (ammonium isocyanate, Demyk et al. [1998;](#page-9-0) Raunier et al. [2003;](#page-10-0) Van Broekhuizen, Keane & Schutte [2004\)](#page-9-0). In each of these cases, the reactants are relatively volatile compounds, likely to sublimate from dust grains at comparatively low temperatures (i.e. *<*100 K for the smallest ones, *<*200 K for the largest ones). On the other hand, the salts formed sublimate at higher temperatures, allowing the snowline of both volatile species involved in the mixtures to be extended towards warmer regions. Therefore, salts are considered good candidates to be involved in the complexation of early chemistry.

In addition, salts would provide an explanation for some of the most significant depletions in the Universe. For example, a nitrogen depletion is observed in several environments including comets (Filacchione et al. [2019\)](#page-10-0) and ISM (Savage & Sembach [1996\)](#page-10-0), compared to the Solar System ratios (Lodders [2010\)](#page-10-0). Likewise, sulphur is reported missing especially in dense cores and protostellar environments (Tieftrunk et al. [1994;](#page-10-0) Goicoechea et al. [2006\)](#page-10-0). The abundance of S-bearing species is then approximately 1000 times lower than that found in diffuse clouds (Kushwahaa et al. [2023\)](#page-10-0), which is also not in tone with their abundance in comets (Calmonte et al. [2016;](#page-9-0) López-Gallifa et al. [2024\)](#page-10-0). So far, there is not one clear explanation to these depletions. Regarding sulphur, it is believed to be trapped in the dust grains cores (Perrero et al. [2024\)](#page-10-0), or in the icy mantles in the form of salts. The latter hypothesis would explain the release of S-containing compounds in comets and enable their detection in larger amounts.

Morevover, the H_2S presence in the early stages of pre-stellar cores development is a determining factor when it comes to understanding the carbonyl sulfide (OCS) as well as the S-bearing molecules content, especially in low-mass stars. The H₂S on interstellar dust grains would indeed be converted to OCS throughout the collapse phase, as long as the core environment is cold enough (typically *<*20 K) to maintain a high quantity of CO ice (Drozdovskaya et al. [2018\)](#page-9-0). The transformation would mostly happen via grain-surface processes(Ferrante et al. [2008;](#page-10-0) El Akel et al. [2022\)](#page-10-0). More elaborately, the rate of UV radiations in the pre-stellar core is also likely to determine the fluency of the conversion of H_2S to OCS, as shown by laboratory experiments where a high UV irradiation involves a high conversion rate (Chen et al. [2015\)](#page-9-0), and observations of IRAS 16293−2422 B, where such a rate would explain the amount and diversity of S-bearing molecules detected compared to the Solar system ratios (Ruffle et al. [1999;](#page-10-0) Jiménez-Escobar, Caro & Chen [2014;](#page-10-0) Martín-Doménech et al. [2016;](#page-10-0) Drozdovskaya et al. [2018;](#page-9-0) Phuong et al. [2018\)](#page-10-0).

It is thus also necessary to take ammonium salts into account in gas-grain models of protostars, as they would significantly influence the proportion of sulphur available to create S-bearing molecules, despite the apparent depletion.

In this paper, we aim to investigate the nature of the products upon the desorption of NH4SH. We propose to densify the laboratory data on the formation, infrared (IR) signatures, and characteristics of the salt on a cold surface representing an interstellar dust grain. We performed experiments and analyses with varying ratios of incoming $NH₃$ and $H₂S$ in pure form, on a bare gold-plated surface, and in the case of a water ice substrate. The reactivity and destruction of the salt under H atoms exposure were also studied. Hereafter, Section 2 describes the experimental methods and setup reproducing the interstellar conditions. In Section [3,](#page-2-0) we present our infrared spectroscopy and mass spectrometry data on ices of $NH₃$ and $H₂S$, on NH4SH salt, on NH4SH exposed to H atoms, and finally in the case of mixtures containing water molecules in different proportions. In Section [4,](#page-7-0) we discuss the astrophysical implications and show that the inclusion of NH4SH in a gas grain model can better explain the OCS/H₂S ratio observed towards protostars.

2 EXPERIMENTAL METHODS

All of the experiments have been performed with the VErs de NoUvelles Synthèses (VENUS) experimental setup, located in LERMA-CY and described in detail elsewhere (Congiu et al. [2020\)](#page-9-0). It consists of a gold-coated circular surface of 9 mm diameter, mounted on a closed-cycle He cryostat and held inside an Ultra-High Vacuum stainless steel chamber. The base pressure is approximately of 10^{-10} mbar. Chemical species under a gaseous, liquid, or solid state at room temperature are kept or converted into gas state in the molecular or atomic beamlines, then sent towards the surface by means of a differential pumping through two intermediate chambers. The beams meet on the surface and the gas-phase interactions above the cold sample are negligible. Indeed, the partial pressure does not exceed 10−⁸ mbar in the volume of the beams, and the main chamber is maintained at 10−¹⁰ mbar even when species are being deposited. The only exception is for H_2 diffusion when the atomic hydrogen beam is being used (the partial pressure of the chamber therefore going up to 10−⁸ mbar). Overall, surface–surface and possibly gas-surface interactions are largely dominant compared to gas–gas interactions. Representing an interstellar dust grain, the surface temperature can vary between 6 and 350 K and is computer-controlled by means of a resistive heater clamped behind the sample holder.

2.1 RAIRS and TPD

A typical experiment consists of two phases. The species are deposited at a constant temperature for a sufficiently long time to allow their detection, forming an ice film on the surface (phase 1, deposition), before being sublimated and released in gaseous form during a linear increase in the temperature of the surface (phase 2, temperature programmed desorption, TPD). During both phases, we can monitor the ice evolution with a Vertex-70 Fourier Transform infrared (FT-IR) spectrometer, performing reflection absorption infrared spectroscopy (RAIRS), a non-destructive analysis technique. The surface is continuously scanned with a grazing infrared beam at an 83◦ angle, over a spectral range of 800–4500 cm−¹ (12.5– $2.22 \,\mu$ m). The light coming from the sample with the desired angle is collected in a N_2 -cooled mercury cadmium telluride (MCT) detector. Every two minutes (corresponding to 256 scans), a mean spectrum is produced.

Once the deposition at low temperature is achieved, the temperature of the surface is linearly increased at a rate $\beta = 0.2 \text{ K s}^{-1}$. The desorbing species are detected with a Hiden 51/3F vertically translatable quadrupole mass spectrometer (QMS), placed 5 mm in front of the surface during the course of the TPD. Every molecule and fragment can then be probed according to their respective mass, chosen by the user and measured in atomic mass units (amu). The ionization energy of the QMS was set at 30 eV, and the number of species entering the QMS head is monitored as a function of temperature. From this, numerous parameters can be derived, including the binding energies between the molecules and the gold surface.

2.2 Atomic, water, and molecular beams

In this work we have carried out a series of experiments, as shown in Table [1,](#page-2-0) for which we provide TPD and infrared data. They can be grouped into several categories: individual species (i, ii),

Exp.	Reactive system	Coverage	$T_{\rm dep}$	H flux $(\text{atoms cm}^{-2} \text{ s}^{-1})$	
		(monolayers)	(K)		
	${H_2S}$	0.1; 0.15; 0.35; 0.65	10		
ii	$\{NH_3\}$	0.35; 0.4; 0.5; 1	10		
iii	${H_2S + H}$	0.65:10	10	8×10^{12}	
iv	${H_2S + NH_3}$	1.8:4.65	10		
\mathbf{V}	${H_2S + NH_3}$	1.8:1.9	10		
vi	${H_2S + NH_3} + {H}$	1.8:1.9:50	80:10	8×10^{12}	
vii	${H_2S + NH_3 + H}$	1.8:1.9:30	10	8×10^{12}	
viii	${H_2S + NH_3 + H_2O}$	1.9:1.9:12	10		
ix	${H_2S + NH_3 + H_2O}$	0.8:1.5:30	10		
\mathbf{x}	$\{NH_3 + H_2O\}$	1:6	10		
xi	${H_2S + NH_3 + H_2O + H}$	1.9:1.9:12:30	10	8×10^{12}	

Table 1. List of the conducted experiments and their associated parameters (composition of the ice, coverage, deposition temperature, and H flux when used).

hydrogenation of H_2S (iii), salt formation (iv, v), salt with H atoms exposure (vi, vii), salt in a water ice matrix (viii, ix), NH_3 with water (x), and salt in a water ice matrix and under H-exposure (xi). When two or more species are deposited on the surface, we denote the codeposition by placing all the species in a pair of curly brackets. For instance, in the figures of this paper, we use the notation {*A* $+ B$ _{doseA:doseB} to refer to a codeposition of species *A* and *B* on the surface, at temperature *T*, with estimated or calculated doses for each species. Conversely, a sequential deposition is indicated by a series of brackets, that is $\{A + B\}^{T_{AB}}_{\text{doseA:}} + \{C\}^{T_C}_{\text{doseC}}.$

The molecular NH₃ beam was calibrated as described in Noble et al. [\(2011\)](#page-10-0). A monolayer (ML), which defines the theoretical filling of all available adsorption sites on the sample and corresponds to ∼ 10¹⁵ molecules cm−2, was reached after 10 and 31 min of deposition, for $NH₃$ and $H₂S$, respectively. This gives corresponding condensation rates – or fluxes – of $\phi_{NH_3} = 1.7 \times 10^{12}$ molecules cm⁻² s⁻¹, and $\phi_{\text{H}_2\text{S}} = 5.4 \times 10^{11}$ molecules cm⁻² s⁻¹. On the other hand, in the case of H_2O , water vapour was leaked into the main chamber via a delivery manifold with known characteristics (a monolayer grown in approximately 5 min, giving $\phi_{\text{H}_2\text{O}} = 3.3 \times 10^{12}$ molecules cm⁻² s⁻¹). For water, the thickness of an ice layer of 1 ML is 3 Å , which is also a good estimate for other species of comparable size.

In the case of hydrogen, only H_2 molecules can be stored in a gas bottle. The atoms were obtained by dissociating the molecules and creating a plasma in a quartz tube surrounded by a microwave cavity (Surfatron). The H atoms have a short residence time on one adsorption site, due to their propensity to diffuse and recombine rapidly to form H_2 . Therefore, their quantity in ML can only be estimated with the maximum number of reactions that can occur in a barrierless reaction such as $NO + H$ (Congiu et al. [2012\)](#page-9-0). One 'monolayer' of H atoms is then considered to be achieved in about 2 min when the flux is of $\phi_H = 8 \times 10^{12}$ atoms cm⁻² s⁻¹. The H₂ dissociation efficiency was measured before each experiment with H atoms, giving an average of ∼62 per cent, with a possible variation of at most 20 per cent between two different experiments.

3 RESULTS

3.1 Single-component ices: NH3 and H2S depositions

A characterization of the individual species was first carried out in order to determine their properties under our experimental conditions. Fig. 1 displays the TPD spectra of 0.6 monolayers of H_2S , and 1 monolayer of NH3. Both species were deposited at 10 K and desorb from the surface at 76 and 90 K respectively, following a first order desorption behaviour. Several fragments can be identified for each

Figure 1. TPD spectra of 0.65 monolayers of H_2S (top panel), and of 1 monolayer of NH_3 (bottom panel). Both species were deposited at $10 K$. Each curve represents a fragment ionized by the QMS. The dotted line stands for $m/z = 16$ when NH₃ is mixed with H₂O.

species, as shown with each curve, giving a unique cracking pattern for each molecule (see upper panel of Fig. [2\)](#page-3-0). The amounts of each fragment depend mainly on the ionization energy of the QMS, which in this work was set at 30 eV . The H₂S and NH₃ major fragments correspond to the ionized unimpaired molecules: $m/z = 34$ for H_2S^+ and $m/z = 17$ for NH_3^+ .

In the case of H_2S , we detect between 3 and 4 per cent of $34S$, an amount that is not considered to significantly affect the nature and kinetics of the reactions with H and $NH₃$. Even when less than 1

Figure 2. Mass spectra of NH_3 (top left), H_2S (top right), and NH_4SH (bottom). The intensities have been normalized to the most abundant fragment for each of the pure species ($m/z = 17$ for NH₃, $m/z = 34$ for H₂S). For NH4SH, the ammonium fragments are normalized to mass 17, while the hydrogen sulfide fragments are normalized to mass 34.

 ML of H_2S is deposited on the sample, the molecules quickly exhibit a multilayer behaviour under our conditions, making it difficult to determine the time after which a monolayer is reached. Moreover, the areas under the main fragment curves of both H_2S and NH_3 are not directly comparable to each other due to their different ionization cross-sections at 30 eV, of respectively $\sigma_{\text{H}_2\text{S}}^{\text{tot}} = 3.18 \text{ Å}^2$ (Vinodkumar et al. [2011\)](#page-10-0) and $\sigma_{NH_3}^{tot} = 1.57 \text{ Å}^2$ (Kumar & Kumar [2020\)](#page-10-0), giving a ratio of $\frac{\sigma_{\rm H_2S}}{\sigma_{\rm NH_3}} \sim 2.02$. Consequently, the determination of the monolayer completion for H_2S (31 min) was based on the value obtained for $NH₃$ (10 min). Measurements in the salt (cf. Section 3.2) further confirm a calculated ratio of ∼1.9 between the two species, which is in good agreement with the total cross-sections found in the literature.

The infrared bands of H_2S and NH_3 are shown in Fig. 3 with asterisks (∗), and the values are listed in Table [2.](#page-4-0) When very little product is present on the surface, the only relevant and usable signature of H₂S is a small absorption band at 2570 cm⁻¹, testifying to the HS[−] stretching mode (*ν*₃), in agreement with other studies (Jiménez-Escobar et al. 2014 ; Loeffler et al. 2015 ; Oba et al. 2018). The detection of H_2S with a grazing infrared beam is rather difficult, due to the very broad nature of its main signature and to the associated low-band strength, determined directly from experimental measurements at $A(S-H) = 1.12 \times 10^{-17}$ cm molecule⁻¹ (Hudson & Gerakines [2018\)](#page-10-0). Ammonia infrared signatures are present at 3380

Figure 3. Infrared spectra of 0.65 monolayers of H₂S (bottom curve) and of 1 monolayer of NH₃ (top curve). The species signatures are shown with \ast , while the diamonds (\Diamond) indicate residual water present on the detector window, that has not been deposited on the surface.

and 1642 cm^{-1} (v_3 and v_4 NH bending modes of NH₃). They also correspond to the known features documented in the literature (Gálvez et al. 2010 ; Kruczkiewicz et al. 2021). On both spectra, as well as in the other spectra presented in this work, the broad, sharp absorption band at 3205 cm⁻¹, indicated by diamonds (\Diamond), is due to residual water condensed on the IR detector window.

3.2 Salt formation: H2S:NH3 mixtures

When mixed together, ammonia and hydrogen sulfide interact to form ammonium hydrosulfide. Fig. [4](#page-4-0) shows the TPD results of experiment (iv) on the left panel, and of experiment (v) on the right panel. In both cases, NH_3 and H_2S gases were codeposited on the surface held at 10 K, via two different molecular beams. In (iv), NH₃ was in excess (4.65 ML) compared to H_2S (1.8 ML), while in (v) both formed an $∼1:1$ mixture, (1.9 ML of NH₃ with 1.8 ML of H_2S). Only the desorption curves of the major fragments resulting from the temperature increase are shown. The absence of signal for the $m/z = 51$ confirms that, being a salt, NH₄SH is an electrovalent compound and does not contain any covalent bond forming a molecule, but rather two ions HS^- and NH_4^+ (Bragin et al. [1977;](#page-9-0) Loeffler et al. [2015\)](#page-10-0). This is the case for many salts (Raunier et al. [2003;](#page-10-0) Noble et al. [2013;](#page-10-0) Kruczkiewicz et al. [2021;](#page-10-0) Gerakines et al. [2024\)](#page-10-0). Therefore, we observe $m/z = 34$ and $m/z = 17$ as the main components of the salt, desorbing at 153 and 154 K in the two experiments. With very little desorption of pure molecular H_2S and $NH₃$, we note that 6 per cent and 18 per cent of $H₂S$ remain pure in the exp. (iv) and (v) respectively, while 17 per cent and 14 per cent of $NH₃$ remain pure in exp. (iv) and (v), respectively. These results, together with those shown before in Fig. 2, confirm the protection of both $NH₃$ and $H₂S$ against the temperature increase. They are stored in salt form at temperatures at which the volatile species would already have desorbed, only to be released in their original form during the sublimation. This behaviour is however not unique to $NH₄SH$; most salts tend to dissociate into two more volatile species, when facing a sufficient temperature increase (Noble et al. [2013;](#page-10-0) Kruczkiewicz et al. [2021;](#page-10-0) Gerakines et al. [2024\)](#page-10-0).

By calculating the ratio between two representative masses of a chemical species, it is possible to evaluate its conservation in the

Table 2. List of infrared features observed in our experiments and found in the literature. ^(a)(Smith [1991\)](#page-10-0), ^(b)(Oba et al. [2018\)](#page-10-0), ^(c)(Kruczkiewicz et al. [2021\)](#page-10-0), (*d*) (Galvez ´ et al. [2010\)](#page-10-0), (*e*) (Ferraro, Sill & Fink [1980\)](#page-10-0), (*^f*) (Bragin et al. [1977\)](#page-9-0), (*g*) (Hudson & Gerakines [2018\)](#page-10-0), (*h*) (Fathe et al. [2006\)](#page-10-0), (*i*) (Loeffler et al. [2015\)](#page-10-0), (*j*) (d'Hedencourt & Allamandola [1986\)](#page-10-0), (*k*) (Bouilloud et al. [2015\)](#page-9-0), (*l*) (Hudgins et al. [1993\)](#page-10-0).

Experiment	Exp. no.		H_2S		NH ₃		NH ₄ SH			H_2O
		v_1 ^(a)	v_3 ^(a,b,g,h)	$v_3(c,d)$	$v_4^{(c,d)}$	NH_4 ⁺ v_4 ^(c,d,i)	v_3 ^(e)	HS^- $v_4^{(e,f)}$	v_1 ^(l)	$v_2^{(j,k)}$
${H_2S}$			2568							
$\{NH_3\}$	ii		$\overline{}$	3383	1642					
${H_2S + NH_3}$	iv		$\overline{}$	3364	1631	1488	3007	1814	-	
	V	$\overline{}$	2567	3371	1631	1483				
${H_2S + NH_3} + {H}$	vi		-	3371	$\overline{}$	1475	3007			
${H_2S + NH_3 + H}$	vii	-	2572	3366	1628	1480	3003			
${H_2S + NH_3 + H_2O}$	viii		2554	3390	1641	1478	-	$\overline{}$	3312	1684
	ix	2527	-	3385	1644	1487	-	$\overline{}$	3336	1681
$\{NH_3 + H_2O\}$	X		$\overline{}$	3398	1643	$\overline{}$			3341	1690
${H_2S + NH_3 + H_2O + H}$	xi		2551	3388	1643	1478	3009	$\qquad \qquad$	3310	1681

Figure 4. TPD spectra of H2S:NH3 mixtures. The left panel shows results of the experiment (iv), where NH3 was more abundant than H2S on the surface. The right panel shows the results of the experiment (v), where NH₃ and H₂S formed a homogeneous mixture on the surface. Only the major fragments are shown. The grey dotted lines indicate the temperatures at which each product desorbs in its pure form.

final detection by the QMS. In fact, with a fixed cracking pattern and non-varying experimental conditions, the ratio between each of the fragments is supposed to remain identical. However, if a reaction involving a different species produces a new product, the latter will have different binding, cohesive, and interaction energies, thus altering the proportions between the masses. To determine whether the salt desorbs as an independent entity (namely a hypothetical molecular NH4SH) or breaks down into two pure species, we calculated the ratios between $m/z = 33$ and 34 for H₂S, and $m/z = 16$ and 17 for NH3. The values are reported in Table 3. We find that the $33/34$ ratio is not different in pure H_2S or in the salt, being constant at ∼0.4. The 16/17 ratio is slightly different when calculated from a TPD of purely NH₃ or of salt (\sim 0.55 and \sim 0.5, respectively). However, the total ratio remains in the same range, indicating the dissociation of the salt into two different species.

In the experiment with an excess of $NH₃$ (iv), there is also a peculiar desorption of $m/z = 17$ between 90 and 142 K, with a maximum at 117 K. It represents 46 per cent of the total ammonia sent on the surface during this experiment. Ammonia ice alone shows a peak at 90 K, but if it is likely to form H-bonds with other components on the surface, especially with water, its desorption is delayed towards higher temperatures, as shown by the dashed grey line of Fig. [1](#page-2-0)

Table 3. Proportions between different *m/z* ratios in individual species and in the salt.

Experiment	33/34	16/17
${H_2S}(i)$	0.41	
$\{NH_3\}$ (ii)		0.56
${H_2S + H}$ (iii)	0.41	
${H2S + NH3}$ – excess of NH ₃ (iv)	0.41	0.51
${H2S + NH3}$ – homog. mixture (v)	0.40	0.48
${H2S + NH3} + {H}$ (vi)	0.41	0.57
${H_2S + NH_3 + H_2O}$ (viii)	0.40	
${H_2S + NH_3 + H_2O}$ (ix)	0.39	

(lower panel) (Suresh et al. [2024\)](#page-10-0). In the case where ammonia is mixed with H_2S , we believe that the same mechanism occurs between NH₃ and the salt. The cracking pattern of ammonia is preserved, suggesting that it is not chemically bonded to any compound, but is only slightly interacting with NH4SH. We refer to these molecules as 'interacting NH₃', as opposed to 'non-interacting NH₃', which characterizes molecules that do not interact with the salt. We also note that some of the 'non-interacting $NH₃$ ' could be related to the

Figure 5. Infrared spectra for the main experiments. *a, c, d, e, f* correspond to experiments iv, v, vi, vii, viii. The spectrum b was recorded at 80 K, during the heating phase after the deposition of ${H₂S + NH₃}$ (iv). The spectrum *d* was recorded at 80 K, before the H bombardment phase of the salt (experiment vi). The arrows ∨ indicate the salt signature and the symbol ∗ denotes the signatures of pure H₂S and NH₃. The symbol \Diamond indicates the bands of H₂O, whether caused by ice condensed on the MCT detector window (*a, b, c, d*), or ice intentionally deposited on the surface (*f*), or condensed on the surface due to residual pressure of H2O in the main chamber (*d*).

uneven coverage between the two beams used to mix NH_3 and H_2S . However, this amount should not exceed 10 per cent.

In the mid-IR region, NH4SH has visible spectral signatures. The spectra associated with simple $H_2S:NH_3$ mixtures (experiments iv, v, and first phase of vi), correspond to the curves *a*, *c*, and *d* in Fig. 5. The absorption bands and corresponding wavenumbers are listed in Table [2.](#page-4-0) Already at 10 K, the most prominent feature observed in all NH4SH experiments performed in this work is the *ν*⁴ NH bending mode of NH₄⁺, which appears around 1485 cm⁻¹. As more NH₄SH is formed when more NH₃ molecules are available on the sample (exp. iv), the absorbance increases and the band becomes more prominent. Similarly, we detect larger NH₃ signatures (v_3 near 3370 cm⁻¹ and v_4 near 1635 cm⁻¹) due to its excess, in the spectrum *a*. Despite a high noise level, mainly caused by the vibrations of the experimental setup and the initial distortion of the baseline, delimited by hatching in Fig. 5, we attribute the 2900–3000 cm⁻¹ absorptions to the NH₄⁺ ion (Ferraro et al. [1980\)](#page-10-0).

As the sample was gradually heated in experiment (iv), we monitored the evolution of the ice with the FT-IR. The main signature of NH₄SH at 1485 cm⁻¹ becomes the sharpest when the temperature of the sample reaches 80 K. This agrees well with the spectrum obtained during the deposition at a constant temperature of 80 K (profile *d*, first phase of exp. vi). Whereas Bragin et al. [\(1977\)](#page-9-0) find that the salt crystallizes around 100 K, Loeffler et al. [\(2015\)](#page-10-0) observed amorphous NH4SH up to 120 K, followed by its crystallization between 130 and 160 K. On the other hand, the crystallization phase of H_2S occurs between 70 and 80 K (Fathe et al. [2006\)](#page-10-0), which corresponds well to the sharp maximum observed under our experimental conditions.

3.3 Salt behaviour under H atoms exposure

To the best of our knowledge, the behaviour of NH4SH exposed to H atoms has not been studied in the laboratory under conditions relevant to the ISM. Studies on NH4SH have mainly focused on

Figure 6. TPD spectra resulting from the $\{NH_4SH\}$ + $\{H\}$ experiment (vi). Only the major fragments are shown. Each of them has been multiplied by 5 for readability and scale consistency purposes. *Inset*: infrared spectra of H2S (bottom curve) and of $H_2S + H$ codeposited (top curve). The grey dashed line shows the absorption band reference (2570 cm^{-1}) .

its behaviour in mixed ices (Altwegg et al. [2020\)](#page-9-0), under photolysis (Lebofsky & Fegley [1976\)](#page-10-0), with temperature changes, under proton irradiation (Loeffler et al. [2015,](#page-10-0) [2016;](#page-10-0) Loeffler & Hudson [2018\)](#page-10-0), and under electron irradiation (Mahjoub et al. [2016\)](#page-10-0). In contrast, the H-exposure of H_2S has been the subject of some experimental works (Oba et al. [2018;](#page-10-0) El Akel et al. [2022;](#page-10-0) Santos, Linnartz & Chuang [2023\)](#page-10-0) and computational studies (Garrod & Herbst [2007;](#page-10-0) Lamberts & Kästner [2017\)](#page-10-0). In this study, three types of H_2S and NH4SH H-exposure experiments were performed (Table [1,](#page-2-0) iii, vi, vii). In all cases, the H atoms flux remained constant at $\phi_H = 8 \times 10^{12}$ atoms cm^{-2} s^{-1}.

We found the exposure of H₂S to H atoms to be very effective. With a continuous flow of hydrogen atoms during the deposition, the final amount of product detected in ${H_2S + H}$ corresponds to about 38 per cent of that of a signal of unaltered ${H_2S}$. However, its infrared signature becomes practically invisible (see inset of Fig. 6), giving usinformation about the signal-to-noise limit for the detection of H2S with our RAIRS instrument, which is about 0.25 ML (2*.*5 × 10¹⁴ molecules cm−2). The loss of product can be explained by the chemical desorption that takes place during the reaction processes

$$
H_2S + H \longrightarrow HS + H_2,\tag{1}
$$

$$
HS + H_2 \longrightarrow H_2S + H. \tag{2}
$$

According to equation (1) and equation (2), once H_2S undergoes H abstraction to form an HS radical, it is likely to either recombine with another H to reform H_2S or undergo chemical desorption, that is to be ejected be ejected from the surface and return to the gas phase due to excess energy that is too large to dissipate (Dulieu et al. [2013\)](#page-9-0). The competition between the different surface processes, together with the $H + H$ addition, has previously been studied quantitatively (Lamberts & Kästner 2017). The cracking patterns of the residual H_2S remain identical, except for the $m/z = 35$ and 36 which are significantly reduced, possibly indicating a preferential chemical desorption of $H_2^{34}S$ and $H^{34}S$ over $H_2^{32}S$.

Exposure of the salt to hydrogen atoms is also effective in reducing the amount of product remaining on the surface, although not as fast as the ${H₂S + H}$ reaction. Fig. [7](#page-6-0) shows the time evolution of the infrared band regions of the main NH4SH signature (at 1485 cm−1) during the experiment (vi). The salt was grown by

Figure 7. Scatter plot of the infrared band areas of the main signature of NH₄SH (1485 cm⁻¹) as a function of time, in the experiment vi. The salt was grown for 60 min at 80 K and then exposed to H atoms for 100 min at 10 K with our atomic beam. The decreasing solid line is a linear fit of the decrease in the IR band area.

one hour ${H_2S + NH_3}$ codeposition on the surface, which was maintained at 80 K for infrared data readability purposes. We were careful to maintain stable conditions to ensure the lowest possible presence of pure, non-interacting $NH₃$ and $H₂S$ on the sample. The surface was then cooled to 10 K before being exposed to hydrogen atoms for 100 min, corresponding to approximately 43 ML of H. The two phases(deposition and H-exposure) are highlighted in Fig. 7 with a blue and red background colour, respectively. We estimate that the salt is destroyed with a half-life of 120 min by sending $\phi_H = 8 \times 10^{12}$ atoms cm⁻² s⁻¹, which is about 1/3 of the destruction obtained after sending ∼ 80 times more hydrogen atoms than the amount of NH₄SH present on the surface. It should be noted, however, that this rate of destruction may depend on the intrinsic mobility of H on the surface, which may be slightly different depending on the experimental conditions. In fact, the reformation of H_2 consumes most of the H atoms sent to the surface. Here, the destruction of the salt can be fitted with a linear rather than an exponential curve, meaning that the reactivity of the salt can be complete or end with a small plateau. We therefore conclude that, while the salt form does not prevent all hydrogenation or abstraction reactions from occurring, it greatly reduces the reactivity rate, compared to when only molecular H_2S is exposed to H.

The TPD traces obtained as a result of this two-phase experiment are shown in the Fig. [6.](#page-5-0) Each fragment curve has been multiplied by 5, in order to keep the same scale as before. Since there is less NH4SH on the surface, the shift of its desorption temperature from 153 to 157 K is not surprising. As mentioned above, the deposition conditions at 80 K were such that the amounts of pure NH_3 and H_2S were as low as possible. Despite this, we detect 1.1 ML of excess NH₃ as its 'interacting' form after H-exposure, in addition to 0.6 ML of NH4 ⁺ constituting the salt, which gives a total amount of 1.7 ML of residual ammonia on the surface (i.e. ∼90 per cent of the total amount of ammonia sent by the beams). An equivalent quantity of 0.55 ML of H2S constituting the salt remains on the sample after sending H atoms (i.e. \sim 31 per cent of the initial H₂S). We also conducted a triple codeposition experiment (vii), in which NH_3 molecules, H_2S molecules, and hydrogen atoms were simultaneously sent to the surface at 10 K, with fluxes identical to experiment (vi). In contrast to the results of experiment (vi), in experiment (vii) small amounts of pure non-interacting NH_3 and H_2S remained on the surface (0.2 ML and 0.1 ML, respectively), in addition to a continuous desorption of $NH₃$ (90–150 K) due to H-bonding (0.85 ML). In this case, between 0.7 and 0.8 ML of each species is incorporated into the salt. Again, it appears that ∼95 per cent of NH3 remains on the surface after Hexposure, whereas for H₂S it is \sim 46 per cent. The infrared spectrum associated with this experiment is shown in Fig. [5\(](#page-5-0)e). Apart from larger NH₃ signatures, no other special features were observed. The small sharp peak at 2342 cm^{-1} corresponds to CO_2 due to a minor leak.

Overall, our results show that the consumption of H_2S when facing a hydrogen atoms exposure, even in salt form, does occur, but at a reduced rate. The reaction of NH4SH with hydrogen atoms may disrupt the ionic crystal structure of the salt to form $NH₃ + SH + H₂$. The process is globally exothermic and is equivalent to the reaction described by reaction [1.](#page-5-0) In this case, the SH formed could react with the excess of hydrogen atoms to reform H_2S , some of which would pass into the gas phase by chemical desorption. This could explain the loss of NH_4SH and H_2S observed experimentally. However, this reaction does not seem possible because the first step should be $NH_4SH + H \longrightarrow NH_3SH + H_2$, which is endothermic and should rather form NH4S. In fact, the ionic crystal energy required to form NH4S is similar to that required to form the salt NH4SH:

$$
NH3 + H2S \longrightarrow NH4+ + SH-(-6.40 eV)
$$
 (3)

$$
NH3 + SH \longrightarrow NH4+ + S-(-6.37 eV)
$$
 (4)

In reaction 4, if there is an excess of hydrogen atoms, NH₄SH should be formed, which then acts as a catalyst for the hydrogen recombination reaction (H + H \rightarrow H₂), without loss of H₂S by chemical desorption.

Another possibility is the formation of the salt NH_4H_2S $(NH_4^+, H_2S, e-)$. With an excess of H, we would then have the following endothermic reaction:

 $NH_4H_2S + H \longrightarrow NH_4^+, H^- + H_2S \longrightarrow NH_3 + H_2S + H_2$ (5)

The general reaction would be:

 $NH_4SH + H + H \longrightarrow NH_3 + H_2S + H_2.$ (6)

Some of the H_2 S formed will give back NH₄SH, but some may also react with hydrogen atoms to form H_2 + SH and lead to a loss of H2S by chemical desorption during its reformation. It is clear from our results that no pure H_2S remains on the surface after 100 min of exposure to H. The fact that $NH₄SH$ is less sensitive to hydrogenation than H_2S can be explained by the fact that a part of the salt is reformed after reaction with H, while the other part is consumed indirectly by the chemical desorption of H_2S according to equation [\(1\)](#page-5-0), and by the fact that the whole mechanism requires several hydrogen atoms.

Looking at the infrared spectra obtained for this experiment, we observe a progressive destruction of the salt signature at 1485 cm−¹ during the H-exposure phase. No particular new features were observed, and no quantification of the NH₃ reformation could be calculated due to a severe baseline deformation.

3.4 Salt behaviour in water ice mixtures

The behaviour of NH4SH in the environment of other compounds was investigated by Altwegg et al. [\(2020\)](#page-9-0), who carried out TPD experiments where $CH₃OH$, $NH₃$, $H₂S$, and $H₂O$ were simultaneously deposited on a gold-coated surface held at 50 K. They observed the formation of ammonium hydrosulfide in the ice mixture, as in previous works where the use of ^{15}N confirmed the NH₃ chemistry

Figure 8. TPD spectra of a H₂S:NH₃:H₂O mixture, where all species have been codeposited at 10 K. The fragment $m/z = 16$ accounts for $[NH_2]^{+}$, to avoid the confusion between $[NH_3]^+$ and $[OH]^+$ that $m/z = 17$ would result in.

in a similar mixture before its irradiation by electrons and thermal processing (Mahjoub et al. [2016\)](#page-10-0).

Other works have studied more generally the behaviour of ammonium salts in water ice mantles. For example, for salts that desorb at temperatures higher than water $(NH_4^+HCOO^-$, $NH_4^+CH3COO^-$, NH_4 ⁺CN⁻, among others), no or little influence of H_2O on the TPD profiles of the salts and on the IR spectra has been registered (Noble et al. [2013;](#page-10-0) Kruczkiewicz et al. [2021\)](#page-10-0). Conversely, for a salt that sublimates in the water desorption region like NH4SH, we found a shift of 10 K in the desorption temperature, as it can be seen in Fig. 8 which shows the results of experiment (viii). The kinetic order of the desorption does not change, but its maximum appears at the very end of the H₂O desorption, towards 163 K. The NH₄SH binding energy would therefore be greater for H_2O than for gold, implying that even in a codeposition situation, NH4SH remains bound to water until there is not enough left to hold it on the now too hot surface. We also note that this happens after the water has crystallized.

Given the double assignment of $m/z = 17$ to ionized OH⁺ and NH_3^+ , $m/z = 16$ has been chosen here to represent NH_3 . In this experiment, there is no non-interacting $NH₃$ on the surface; only 0.6 ML of ammonia interacts with either water or NH4SH. In comparison, 70 per cent and 86 per cent of the total quantity of ammonia and hydrogen sulfide are constituting the NH4SH salt.

Regarding the infrared results, the major changes undoubtedly concern the water band, located in the absorption regions of NH3 $(3400 \text{ cm}^{-1}$ for the −OH stretching mode and 1650 cm⁻¹ for the H2O bending mode). Additionally, a slight blue-shift occurs for the H₂S absorption, measured at 2554 cm⁻¹ instead of \sim 2570 cm⁻¹ (see Table [2\)](#page-4-0).

In order to verify the influence of a water ice matrix on the desorption of the salt under more astronomically relevant conditions, we performed experiment (ix). The TPD results are displayed in Fig. 9. The IR spectrum is not shown, due to the lack of signal other than H₂O. Approximately 0.8 ML of H₂S and 1.5 ML of NH₃ have been diluted in 30 ML of $H₂O$ at 10 K, as a result of regular valve openings throughout the deposition. The readability of the results is therefore improved by using a logarithmic scale and dividing the H2O component by a factor of 10. As in experiment (viii), very small returns from NH₄SH to H₂S and NH₃ are visible (0.1 ML of each, approximately). The desorption is also shifted to 163 K, following the same pattern as the end of the water desorption. Our results are

Figure 9. TPD spectra of a H₂S:NH₃:H₂O mixture in more astrnomically relevant conditions. About 1.5 ML of NH₃ and 0.8 ML of H₂S have been deposited bit by bit in a water ice matrix of ∼30 ML. A logarithmic scale is used for readability purposes.

consistent with the single NH4SH TPD spectrum published to date, where the desorption peak occurs at 168 K in a mixed ice containing H_2O , CH₃OH, H₂S, and NH₃ (Altwegg et al. [2020\)](#page-9-0).

Overall, water seems to affect the desorption and decomposition of NH4SH only by delaying it, shifting its temperature by 10 K under our experimental conditions, regardless of the mixture ratios between different species.

4 ASTROPHYSICAL IMPLICATIONS

H2S and OCS are the main sulphur species observed in the gas phase in protostars (Kushwahaa et al. [2023\)](#page-10-0), with variable ratios. It is therefore important to understand the evolution of their ratios through modelling, as well as the potential role of salts that form when *s-*H2S (*s-* meaning solid, on surfaces) is mobile on ice, during a temperature increase. To model the abundances in protostars, we used the Nautilus gas–grain model in its three-phase form (Ruaud, Wakelam & Hersant [2016\)](#page-10-0). The model was used to simulate the abundances of atoms and molecules in neutral and ionic form as a function of time, employing kida.uva.2014 (Wakelam et al. [2015\)](#page-10-0). The basic chemical network has recently been updated to better describe COMs on grains and in the gas phase (Manigand et al. 2021 ; Loison et al. 2022 ; Agúndez et al. 2023), but also to describe the formation of salts. 800 species are identified in the network and are involved in 9000 individual reactions. In the model that we used, the elements are initially in their atomic or ionic form (the ones with an ionization potential below 13.6 eV are considered to be fully ionized), and the C/O elemental ratio is equal to 0.71. The initial sulphur abundance is set to 6.0×10^{-7} with respect to H, corresponding to a depletion by a factor of 20 compared to the typical cosmic abundance in dense molecular clouds (Fuente et al. [2016\)](#page-10-0). The grain surface and the mantle are both chemically active in these simulations, while the accretion and desorption processes are allowed only between the surface and the gas phase, except for the sputtering by cosmic rays. The dust-to-gas ratio is 0.01 by mass. A sticking probability of 1 is assumed for all neutral species, and desorption can occur via thermal and non-thermal processes, including chemical desorption and sputtering of ices by cosmic ray collisions (Wakelam et al. [2021\)](#page-10-0). The formalism of the surface reactions and a more

Figure 10. Fractional abundances (n/n_{H_2}) of H₂S, and NH₄SH (left panels) and OCS (right panels) in gas phase (black lines) and on the grain surface (blue and purple lines) as a function of time during the collapse phase predicted by the Nautilus astrochemical model. The time axis is reversed to better visualize the abundance evolution. Upper panels: without NH₄SH formation pathway. Lower panels: with inclusion of NH₄SH formation pathway. The abundance of sulphur initially in the gas phase is fixed at ¹*.*² [×] ¹⁰−6, that is ^a depletion factor of ²⁰ compared to the solar abundance (Asplund et al. [2009\)](#page-9-0). The grey rectangle at the top right of each pannel represents the observed abundances towards IRAS 16293−2422B (Kushwahaa et al. [2023\)](#page-10-0), assuming an uncertainty of 3 \times 10⁻⁷ in the *y-*axis.

detailed description of the simulations can be found in Ruaud et al. [\(2016\)](#page-10-0).

To study the chemistry in protostellar envelopes, we have used the same approach as in Manigand et al. [\(2021\)](#page-10-0). We perform two steps: a prestellar and a collapse phase. In the first step, conditions typical of the cold core are assumed for 10^6 yr (gas and dust temperature of 10 K, a proton density of 2×10^4 cm⁻³, a visual extinction of 4.5 mag, a cosmic ray ionization rate of 1.3×10^{-17} s⁻¹, and a standard external UV field of 1 G0). The final composition of the cold core is then used as the initial condition for the collapse phase, which is the same as in Aikawa et al. [\(2008,](#page-9-0) [2012\)](#page-9-0), Wakelam et al. [\(2014\)](#page-10-0), and Manigand et al. [\(2021\)](#page-10-0). The selected parcel of gas begins to collapse at a radius of about 10 000 au and ends-up at a radius of 15 au, where the gas temperature reaches ∼265 K and a density of 4×10^8 cm⁻³.

In the version of the NAUTILUS code used in this work, we have introduced two types of salts: salts stable in the gas phase and ammonium salts, that would form by the Langmuir–Hinshelwood ered as classical molecules, even if they dissociate as ions in ice (Acharyya, Woon & Herbst [2023\)](#page-9-0). On the other hand, ammonium salts do not exist in the gas phase and are only stable when the ions resulting from a proton transfer from an acid to the base (here, ammonia) are sufficiently stabilized by interactions with water molecules, in order to compensate for the exothermicity of the proton transfer. Calculating the interaction energy between ions and the structure of amorphous ices is therefore very complicated, as it is necessary to estimate for which systems these energies are sufficient to ensure the formation of ion pairs, that is salts. To do this, we have used a method similar to that of Woon [\(2012\)](#page-10-0), which consists in calculating the most stable structures for acid-base (H2O)*ⁿ* systems by varying the size of the water cluster. If a proton transfer occurs for $n \leq 8$, we consider the salt formation, barrierless when exothermic (Woon [2012\)](#page-10-0). The ammonium salts thus introduced are NH_4 ⁺CN⁻, NH₄⁺Cl⁻, NH₄⁺NCO⁻, NH₄⁺NCS⁻, NH₄⁺HCOO⁻, and NH_4 ⁺SH⁻. This would explain, among other things, the ice

mechanism. The former, such as table salt Na⁺Cl[−], can be consid-

absorption band usually observed at 2165 cm^{-1} (4.62 µm) and attributed to OCN−. However, since the bands attributed to OCN[−] and NH4 ⁺ are seen in dense molecular cloud ices (McClure et al. [2023\)](#page-10-0), it also implies that either NH3, HNCO, or HOCN are mobile at low temperatures, so that the diffusion energy is much lower than the usual value of 0.4 times the adsorption energy.

In this first version of the NAUTILUS code with salts taken into account, and despite our experimental results showing a significant reactivity, even reduced, of salts with hydrogen atoms, we do not include salt destruction reactions except for desorption. Due to the instability of ammonium salts in the gas phase, unlike usual molecules, we cannot calculate their desorption temperature with the usual models, including the proxy approach developed in Wakelam et al. [\(2017\)](#page-10-0). We have therefore chosen to set the desorption energies for the destructive process s -NH₄SH \rightarrow NH₃+ H₂S so as to reproduce the experimental sublimation temperature of about 150 K for *s*-NH4SH, corresponding to a binding energy of about 6000 K, slightly larger than that of water.

In models of ice formation in interstellar clouds, the gas-phase reactions of $S⁺$ and S are not efficient enough to consume significant amounts of sulphur, which will eventually stick to the ice. The evolution of solid sulphur at low temperatures in dense clouds (around 10 K) is then dominated by reactions with atomic hydrogen, with the production of s -H₂S. As the H–SH bond of H₂S is weaker than the H–H bond of H2, *s-*H2S tends to react with *s-*H by tunneling effect, giving $s-H_2$ + $s-SH$. The $s-SH$ radical can also react with *s-*H to give mainly *s-*H2S (and probably *s-*S), creating a loop and thus an equilibrium between *s-*H2S and *s-*SH (and *s-*S) (Lamberts & Kästner [2017;](#page-10-0) Vidal et al. 2017; Oba et al. [2018\)](#page-10-0). Since both *s*-SH and *s-*H2S have similar abundances at 10 K in our model and are the main reservoir of sulphur for evolved dense clouds, this equilibrium is controlled by the abundance and mobility of H on the gains.

The resulting model predictions are shown in Fig. [10.](#page-8-0) Inclusion of the experimental results in the NAUTILUS astrochemical gas–grain model corroborates the importance of solid-state NH4SH to better constrain the H_2S and OCS abundances (Fig. [10,](#page-8-0) lower panels). When the dense cloud gives rise to the formation of protostars, the temperature in the collapse phase increases, leading to a significant increase in the mobilities of the various species present on the grains, especially radicals. Among the radicals we count *s-*SH and *s-*S, but also *s-*HCO, i.e. the most abundant radical on the grains due to that of *s*-CO, together with notable abundances of *s*-CH₂OH and *s*-CH₃O. In our model, the mobilities of *s-*SH, *s-*S, and *s-*HCO are comparable and inferred from the adsorption energies of Wakelam et al. [\(2017\)](#page-10-0). The reaction of *s-*S with *s-*HCO gives *s-*SH + *s-*CO as well as *s-*H + *s-*OCS, while the reaction of *s-*SH and *s-*HCO can give *s-*HCOSH, $s-H_2S + s-CO$, and $s-H_2 + s-CS$. The solid $s-HCOSH$ species reacts quite easily with H by tunneling effect and are converted fairly rapidly to *s-*OCS. HCOSH was detected towards quiescent clouds in the Galactic Centre (Rodríguez-Almeida et al. [2021;](#page-10-0) García de la Concepción et al. [2022\)](#page-10-0), which is in good agreement with our model, where the abundance of *s-*HCOSH in protostars is low.

5 CONCLUSIONS

In our experimental study, we took a closer look at the formation, reactivity, and desorption of ammonium hydrosulfide salts. These salts are in fact chemical intermediates with specific properties that give them an important role in molecular evolution during star formation processes, but also in the sublimation of ices in general, as in comets. In particular, we have shown that:

(i) NH4SH salts are formed at temperatures as low 10 K from a mixture of H_2S and NH₃, and appear to crystallize around 80 K under our experimental conditions.

(ii) Even diluted in water, H_2S , and NH_3 mixtures lead to the formation of ammonium hydrosulfide salts, and their desorption is shifted to 163 K, that is 10 K higher than in the $H_2S:NH_3$ case, regardless of the mixture ratios in the $H_2O:H_2S:NH_3$ system.

(iii) In astrophysical contexts (i.e. protostars or protostellar discs), salts are likely to desorb with water, and in any case at higher temperatures than the individual pure species NH_3 and H_2S .

(iv) Adopting an effective binding energy of ∼6000 K for NH4SH and the release of the two components NH_3 and H_2S in the gas phase seems a reasonable description of the desorption process.

(v) Including salts in present chemical networks helps to find a better agreement with the observed OCS/H₂S ratio in protostars.

(vi) S-bearing ionic compounds could explain the non-detection of H2S in ices to date, as some of it would be in the form of salts.

(vii) Salts do not provide total protection of H_2S against exposure to hydrogen atoms, but they do slow down the H-abstraction process.

ACKNOWLEDGEMENTS

This work was funded by CY Initiative of Excellence (grant 'Investissements d'Avenir' ANR-16-IDEX-0008), Agence Nationale de la recherche (ANR) SIRC project (Grant ANR-SPV2024482020– 2024), by the Programme National 'Planétologie' and 'Physique et Chimie du Milieu Interstellaire' (PCMI) of CNRS/INSU with INC/INP co-funded by CEA and CNES.

DATA AVA IL AB IL IT Y

The experimental data and conditions have been saved on our own laboratory server, under the form of metadata files associated with each experiment. They will be released publicly after a minimum of 3 yr after their production. For the time being, the data are available from the corresponding authors upon reasonable request.

REFERENCES

- [Acharyya](#page-8-0) [K.,](#page-8-0) Woon D. E., Herbst E., 2023, [MNRAS,](http://dx.doi.org/10.1093/mnras/stad3262) 527, 1722
- Agúndez [M.](#page-7-0) et al., 2023, [A&A,](http://dx.doi.org/10.1051/0004-6361/202346076) 673, A34
- [Aikawa](#page-8-0) [Y.,](#page-8-0) Wakelam V., Garrod R. T., Herbst E., 2008, [ApJ,](http://dx.doi.org/10.1086/524096) 674, 984
- [Aikawa](#page-8-0) [Y.,](#page-8-0) Wakelam V., Hersant F., Garrod R. T., Herbst E., 2012, [ApJ,](http://dx.doi.org/10.1088/0004-637X/760/1/40) 760, 40
- [Altwegg](#page-0-0) [K.](#page-0-0) et al., 2020, Nat. [Astron.,](http://dx.doi.org/10.1038/s41550-019-0991-9) 4, 533
- Asplund M., Grevesse N., Sauval A. J., Scott P., 2009, [ARA&A,](http://dx.doi.org/doi:10.1146/annurev.astro.46.060407.145222) 47, 481
- [Atreya](#page-0-0) S. [K.,](#page-0-0) Wong M. H., Owen T. C., Mahaffy P. R., Niemann H. B., Pater I. D., Drossart P., Encrenaz T., 1999, [Planet.](http://dx.doi.org/10.1016/s0032-0633(99)00047-1) Space Sci., 47, 1243
- [Bergner](#page-0-0) J. [B.,](#page-0-0) Öberg K. I., Rajappan M., Fayolle E. C., 2016, [ApJ,](http://dx.doi.org/10.3847/0004-637x/829/2/85) 829, 85 [Bouilloud](#page-4-0) [M.,](#page-4-0) Fray N., Bénilan Y., Cottin H., Gazeau M.-C., Jolly A., 2015,
- [MNRAS,](http://dx.doi.org/10.1093/mnras/stv1021) 451, 2145 [Bragin](#page-4-0) [J.,](#page-4-0) Diem M., Guthals D., Chang S., 1977, J. [Chem.](http://dx.doi.org/10.1063/1.434936) Phys., 67, 1247
- Van [Broekhuizen](#page-1-0) F. [A.,](#page-1-0) Keane J. V., Schutte W. A., 2004, [A&A,](http://dx.doi.org/10.1051/0004-6361:20034161) 415, 425
- [Calmonte](#page-1-0) [U.](#page-1-0) et al., 2016, [MNRAS,](http://dx.doi.org/10.1093/mnras/stw2601) 462, S253
- [Chen](#page-1-0) [Y.](#page-1-0) J. et al., 2015, [ApJ,](http://dx.doi.org/10.1088/0004-637X/798/2/80) 798, 80
- [Clutter](#page-0-0) D. [R.,](#page-0-0) Thompson W. E., 1969, J. [Chem.](http://dx.doi.org/10.1063/1.1671702) Phys., 51, 153
- [Congiu](#page-2-0) [E.](#page-2-0) et al., 2012, [ApJ,](http://dx.doi.org/10.1088/2041-8205/750/1/L12) 750, L12
- [Congiu](#page-1-0) [E.,](#page-1-0) Sow A., Nguyen T., Baouche S., Dulieu F., 2020, Rev. Sci. [Instrum.,](http://dx.doi.org/10.1063/5.0018926) 91, 124504
- [Demyk](#page-1-0) [K.,](#page-1-0) Dartois E., D'Hendecourt L., Muizon M. J. D., Heras A. M., Breitfellner M., 1998, A&A, 339, 553
- [Drozdovskaya](#page-1-0) [M.](#page-1-0) N. et al., 2018, [MNRAS,](http://dx.doi.org/10.1093/mnras/sty462) 476, 4949
- [Dulieu](#page-5-0) [F.,](#page-5-0) Congiu E., Noble J., Baouche S., Chaabouni H., Moudens A., Minissale M., Cazaux S., 2013, Sci. [Rep.,](http://dx.doi.org/10.1038/srep01338) 3, 1338
- El [Akel](#page-1-0) [M.,](#page-1-0) Kristensen L. E., Gal R. L., Van Der Walt S. J., Pitts R. L., Dulieu F., 2022, [A&A,](http://dx.doi.org/10.1051/0004-6361/202141810) 659, A100
- [Fathe](#page-4-0) [K.,](#page-4-0) Holt J. S., Oxley S. P., Pursell C. J., 2006, J. Phys. [Chem.](http://dx.doi.org/10.1021/jp0634104) A, 110, 10793
- [Ferrante](#page-1-0) R. [H.,](#page-1-0) Moore M. H., Spiliotis M. M., Hudson R. L., 2008, [ApJ,](http://dx.doi.org/10.1086/590362) 684, 1210
- [Ferraro](#page-4-0) J. [R.,](#page-4-0) Sill G., Fink U., 1980, Appl. [Spectrosc.,](http://dx.doi.org/10.1366/0003702804731339) 34, 525
- [Filacchione](#page-1-0) [G.](#page-1-0) et al., 2019, [Space](http://dx.doi.org/10.1007/s11214-019-0580-3) Sci. Rev., 215, 19
- [Friedson](#page-0-0) [A.](#page-0-0) J., 2005, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2005.03.004) 177, 1
- [Fuente](#page-7-0) [A.](#page-7-0) et al., 2016, [A&A,](http://dx.doi.org/10.1051/0004-6361/201628285) 593, A94
- Gálvez [O.,](#page-4-0) Maté B., Herrero V. J., Escribano R., 2010, [ApJ,](http://dx.doi.org/10.1088/0004-637X/724/1/539) 724, 539
- García de la Concepción [J.](#page-9-0) et al., 2022, [A&A,](http://dx.doi.org/10.1051/0004-6361/202142287) 658, A150
- [Garrod](#page-5-0) [R.](#page-5-0) T., Herbst E., 2007, [A&A,](http://dx.doi.org/10.1051/0004-6361:20066704) 467, 1103
- [Gerakines](#page-1-0) P. [A.,](#page-1-0) Moore M. H., Hudson R. L., 2004, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2004.02.005) 170, 202
- [Gerakines](#page-1-0) P. [A.,](#page-1-0) Yarnall Y. Y., Hudson R. L., 2024, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2024.116007) 413, 116007
- [Goicoechea](#page-1-0) J. [R.,](#page-1-0) Pety J., Gerin M., Teyssier D., Roueff E., Hily-Blant P., Baek S., 2006, [A&A,](http://dx.doi.org/10.1051/0004-6361:20065260) 456, 565
- [d'Hedencourt](#page-4-0) L. [B.,](#page-4-0) Allamandola L. J., 1986, A&AS, 64, 453
- [Hudgins](#page-4-0) D. [M.,](#page-4-0) Sandford S. A., Allamandola L. J., Tielens A. G. G. M., 1993, ApJS, 86, 713
- [Hudson](#page-4-0) R. [L.,](#page-4-0) Gerakines P. A., 2018, [ApJ,](http://dx.doi.org/10.3847/1538-4357/aae52a) 867, 138
- Jiménez-Escobar [A.,](#page-1-0) Caro G. M. M., Chen Y. J., 2014, [MNRAS,](http://dx.doi.org/10.1093/mnras/stu1100) 443, 343
- [Kruczkiewicz](#page-0-0) [F.,](#page-0-0) Vitorino J., Congiu E., Theulé P., Dulieu F., 2021, [A&A,](http://dx.doi.org/10.1051/0004) 686, A236
- [Kumar](#page-3-0) [Y.,](#page-3-0) Kumar M., 2020, [Chem.](http://dx.doi.org/10.1016/j.cplett.2019.137071) Phys. Lett., 740, 137071
- [Kushwahaa](#page-1-0) [T.,](#page-1-0) Drozdovskaya M. N., Tychoniec L., Tabone B., 2023, [A&A,](http://dx.doi.org/10.1051/0004-6361/202245097) 672, A122
- [Lamberts](#page-5-0) [T.,](#page-5-0) Kästner J., 2017, J. Phys. [Chem.](http://dx.doi.org/10.1021/acs.jpca.7b10296) A, 121, 9736
- [Lebofsky](#page-0-0) [L.,](#page-0-0) Fegley M., 1976, [Icarus,](http://dx.doi.org/10.1016/0019-1035(76)90151-2) 28, 379
- [Li](#page-0-0) [C.](#page-0-0) et al., 2020, Nat. [Astron.,](http://dx.doi.org/10.1038/s41550-020-1009-3) 4, 609
- [Lodders](#page-1-0) [K.,](#page-1-0) 2010, Solar System Abundances of the Elements, Astrophysics and Space Science Proceedings. Springer-Verlag, Berlin
- [Loeffler](#page-0-0) [M.](#page-0-0) J., Hudson R. L., 2018, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2017.10.041) 302, 418
- [Loeffler](#page-0-0) [M.](#page-0-0) J., Hudson R. L., Chanover N. J., Simon A. A., 2015, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2015.06.015) 258, 181
- [Loeffler](#page-0-0) [M.](#page-0-0) J., Hudson R. L., Chanover N. J., Simon A. A., 2016, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2016.02.010) 271, 265
- [Loison](#page-7-0) A. C. [J.-C.,](#page-7-0) Boulanger A., Caux E., Müller H. S. P., Wakelam V., Manigand S., Jørgensen J. K., 2022, [A&A,](http://dx.doi.org/10.1051/0004-6361/202243038) 660, L6
- López-Gallifa \acute{A} . et al., 2024, MNRAS, 529, 3244
- [Mahjoub](#page-5-0) [A.](#page-5-0) et al., 2016, [ApJ,](http://dx.doi.org/10.3847/0004-637x/820/2/141) 820, 141
- [Manigand](#page-7-0) [S.](#page-7-0) et al., 2021, [A&A,](http://dx.doi.org/10.1051/0004-6361/202038113) 645, A53
- Martín-Doménech [R.,](#page-1-0) Jiménez-Serra I., Caro G. M. M., Müller H. S., Occhiogrosso A., Testi L., Woods P. M., Viti S., 2016, [A&A,](http://dx.doi.org/10.1051/0004-6361/201526271) 585, A112
- [McClure](#page-9-0) [M.](#page-9-0) K. et al., 2023, Nat. [Astron.,](http://dx.doi.org/10.1038/s41550-022-01875-w) 7, 431
- [Niemann](#page-0-0) [H.](#page-0-0) B. et al., 1998, J. [Geophys.](http://dx.doi.org/10.1029/98JE01050) Res.: Planets, 103, 22831
- [Noble](#page-2-0) J. [A.,](#page-2-0) Dulieu F., Congiu E., Fraser H. J., 2011, [ApJ,](http://dx.doi.org/10.1088/0004-637X/735/2/121) 735, 121
- [Noble](#page-1-0) J. [A.,](#page-1-0) Theule P., Borget F., Danger G., Chomat M., Duvernay F., Mispelaer F., Chiavassa T., 2013, [MNRAS,](http://dx.doi.org/10.1093/mnras/sts272) 428, 3262
- [Oba](#page-4-0) [Y.,](#page-4-0) Tomaru T., Lamberts T., Kouchi A., Watanabe N., 2018, Nat. [Astron.,](http://dx.doi.org/10.1038/s41550-018-0380-9) 2, 228
- [Pater](#page-0-0) I. [D.,](#page-0-0) Dunn D., Romani P., Zahnle K., 2001, [Icarus,](http://dx.doi.org/10.1006/icar.2000.6527) 149, 66
- [Perrero](#page-1-0) [J.,](#page-1-0) Beitia-Antero L., Fuente A., Ugliengo P., Rimola A., 2024, [MNRAS,](http://dx.doi.org/10.1093/mnras/stad3896) 527, 10697
- [Phuong](#page-1-0) [N.](#page-1-0) et al., 2018, [A&A,](http://dx.doi.org/10.1051/0004-6361/201833766) 616, L5
- [Raunier](#page-1-0) [S.,](#page-1-0) [Chiavassa](http://dx.doi.org/10.1016/S0009-2614(02)01919-X) T., Marinelli F., Allouche A., Aycard J. P., 2003, Chem. Phys. Lett., 368, 594
- Rodríguez-Almeida [L.](#page-9-0) F. et al., 2021, [ApJ,](http://dx.doi.org/10.3847/2041-8213/abf7cb) 912, L11
- [Roman](#page-0-0) [M.](#page-0-0) T., Banfield D., Gierasch P. J., 2013, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2013.03.015) 225, 93
- [Ruaud](#page-7-0) [M.,](#page-7-0) Wakelam V., Hersant F., 2016, [MNRAS,](http://dx.doi.org/10.1093/mnras/stw887) 459, 3756
- [Ruffle](#page-1-0) [D.](#page-1-0) P., Hartquist T. W., Caselli P., Williams D. A., 1999, [MNRAS,](http://dx.doi.org/10.1046/j.1365-8711.1999.02562.x) 306, 691
- [Santos](#page-5-0) J. [C.,](#page-5-0) Linnartz H., Chuang K.-J., 2023, A&A, 678, A112
- [Savage](#page-1-0) B. [D.,](#page-1-0) Sembach K. R., 1996, [ApJ,](http://dx.doi.org/10.1086/177919) 470, 893
- [Smith](#page-4-0) R. [G.,](#page-4-0) 1991, [MNRAS,](http://dx.doi.org/10.1093/mnras/249.1.172) 249, 172
- [Suresh](#page-4-0) S. [K.,](#page-4-0) Dulieu F., Vitorino J., Caselli P., 2024, [A&A,](http://dx.doi.org/10.1051/0004-6361/202245775) 682, A163
- Theulé [P.,](#page-0-0) Duvernay F., Danger G., Borget F., Bossa J. B., Vinogradoff V., Mispelaer F., Chiavassa T., 2013, Adv. [Space](http://dx.doi.org/10.1016/j.asr.2013.06.034) Res., 52, 1567
- [Tieftrunk](#page-1-0) [A.,](#page-1-0) des Forets G. P., Schilke P., Walmsley C. M., 1994, A&A, 289, 579
- [Vidal](#page-9-0) T. H. [G.,](#page-9-0) Loison J.-C., Jaziri A. Y., Ruaud M., Gratier P., Wakelam V., 2017, [MNRAS,](http://dx.doi.org/10.1093/mnras/stx828) 469, 435
- [Vinodkumar](#page-3-0) [M.,](#page-3-0) Bhutadia H., [Limbachiya](http://dx.doi.org/10.1016/j.ijms.2011.07.017) C., Joshipura K. N., 2011, Int. J. Mass Spectrom., 308, 35
- [Wakelam](#page-8-0) [V.,](#page-8-0) Vastel C., Aikawa Y., Coutens A., Bottinelli S., Caux E., 2014, [MNRAS,](http://dx.doi.org/10.1093/mnras/stu1920) 445, 2854
- [Wakelam](#page-7-0) [V.](#page-7-0) et al., 2015, [ApJS,](http://dx.doi.org/10.1088/0067-0049/217/2/20) 217, 20
- [Wakelam](#page-9-0) [V.,](#page-9-0) Loison J. C., Mereau R., Ruaud M., 2017, Mol. [Astrophys.,](http://dx.doi.org/10.1016/j.molap.2017.01.002) 6, 22
- [Wakelam](#page-7-0) [V.,](#page-7-0) Dartois E., Chabot M., Spezzano S., Navarro-Almaida D., Loison J.-C., Fuente A., 2021, [A&A,](http://dx.doi.org/10.1051/0004-6361/202039855) 652, A63
- [Weidenschilling](#page-0-0) [S.](#page-0-0) J., Lewis J. S., 1973, [Icarus,](http://dx.doi.org/10.1016/0019-1035(73)90019-5) 20, 465
- [Wong](#page-0-0) [M.](#page-0-0) H., Mahaffy P. R., Atreya S. K., Niemann H. B., Owen T. C., 2004, [Icarus,](http://dx.doi.org/10.1016/j.icarus.2004.04.010) 171, 153
- [Woon](#page-8-0) [D.](#page-8-0) E., 2012, [Comput.](http://dx.doi.org/10.1016/j.comptc.2012.01.018) Theor. Chem., 984, 108

This paper has been typeset from a TEX/LATEX file prepared by the author.

© 2024 The Author(s). Published by Oxford University Press on behalf of Royal Astronomical Society. This is an Open Access article distributed under the terms of the Creative Commons Attribution License [\(https://creativecommons.org/licenses/by/4.0/\)](https://creativecommons.org/licenses/by/4.0/), which permits unrestricted reuse, distribution, and reproduction in any medium, provided the original work is properly cited.